## On the Thermodynamic Consistency of the Temperature Dependence of the Debye Temperature

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If the Debye temperature  $(\Theta)$  has temperature dependence, then the expression for lattice heat capacity at constant volume is

$$C^* = C/3 N_A n k_b = (C_D/3 N_A n k_b) [1 \tilde{n} (T/\Theta)(d\Theta/dT)_v]^2 \tilde{n}$$

$$\tilde{n} [ (3/8) + (T/\Theta) D(\Theta/T) ] T (d^2 \Theta/d T^2)_v.$$
 (1)

Here,  $N_A$  is the Avogadro number, n is the number of ions in molecule,  $k_b$  is the Boltzmann constant, D(x) is the Debye function, which have form

$$D(x) = (3 / x^3) \int_{0}^{X} \{ t^3 / [exp(t) \tilde{n} 1] \} dt,$$

$$C_D / 3 N_A n k_b = 4 D(\Theta / T) \tilde{n} 3(\Theta / T) / [exp(\Theta / T) \tilde{n} 1] .$$

When the temperature tends to zero the function (1) must smoothly decrease to zero (it is the third law of thermodynamics). This condition requires that the dependence of  $\Theta(T)$  satisfy the demands: the low-temperature branch of  $\Theta(T)$  at  $0 \le T \le \Theta$  must vary with T as

$$\Theta(T)_{low} \cong \Theta_0 \left[ 1 \, \tilde{n} \, \chi \, (T / \Theta_0)^k \right] \quad , \quad \text{where:} \quad \Theta_0 = \frac{lim \, \Theta(T)}{T \to 0 \, K} \quad ; \quad \chi \ge 0 \; ; \quad k \ge 4 \; . \quad (2)$$

At high temperature, the function (1) must increase, tending at  $T/\Theta \to \infty$  to dependence

$$C^*_{\text{high}} \cong [1 \ \tilde{n} \ (1/20)(\Theta \ / \ T \ )^2] \ [1 \ \tilde{n} \ (\ T \ / \ \Theta)(d \ \Theta \ / \ d \ T \ )_v \ ]^2 \ \tilde{n} \ (\ T^2 \ / \ \Theta)(d^2 \ \Theta \ / \ d \ T^2 \ )_v \ .$$

This leads to the next form of high temperature branch of  $\Theta(T)$ 

$$\Theta(T)_{high} \cong \Theta_{\infty} \exp(\tilde{n} \alpha \Theta_{\infty} / T), \quad \text{where:} \quad \Theta_{\infty} = \frac{\lim \Theta(T)}{T / \Theta_{\infty} \to \infty}; \quad \alpha \ge 0. \quad (3)$$

Value  $\alpha$  is a fitting parameter. The unification of the low temperature branch (2) with the high temperature branch (3) leads to the function  $\Theta(T)$  that can be employed at all temperature intervals

$$\Theta(T) = \Theta_0 \exp \left[ \tilde{n} \chi \left( T / \Theta_0 \right)^k \right] + \Theta_\infty \exp \left( \tilde{n} \alpha \Theta_\infty / T \right)$$
 (4)

Thus, the function  $\Theta(T)$  smoothly decreases from  $\Theta_0$  (at T=0 K) to a minimum, and then it smoothly increases, tending to an asymptotic value  $\Theta_\infty$  at  $T/\Theta \to \infty$ . For calculation of function  $\Theta(T)$  it is necessary to define of the five parameters:  $\Theta_0$ ,  $\chi$ , k,  $\Theta_\infty$ ,  $\alpha$ . In this way, function  $\Theta(T)$  doesn't lead to results that contradict the laws of thermodynamics.

The values of  $\Theta_0$  and  $\Theta_\infty$  are defined for many materials and are in handbooks. A method for the definition of other parameters: $\chi$ , k,  $\alpha$  by means of experimental data is proposed. It is shown that for simple monatomic materials at  $T\approx 0$  K,  $k\cong 4$ .